

Computation of heat asymptotics for nonminimal Laplace type operators: from theory to computer

(Work in progress with Bruno Iochem)

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The problem

- V a smooth hermitean vector bundle V (fiber \mathbb{C}^N) over a compact d -dimensional boundaryless Riemannian manifold (M, g) .
- P a nonminimal Laplace type operator acting on smooth sections $\Gamma(V)$:

$$\begin{aligned} P &= -[g^{\mu\nu} u \partial_\mu \partial_\nu + v^\mu \partial_\nu + w] \\ &= -g^{\mu\nu} u \nabla_\mu \nabla_\nu - (p^\nu + g^{\mu\nu} (\nabla_\mu u) - \Gamma^\nu u) \nabla_\nu - q \end{aligned}$$

- ▶ ∇_μ gauge (only) covariant derivative,
- ▶ u, v^μ, w, p^μ, q are $M_N(\mathbb{C})$ matrix valued functions,
- ▶ $u(x)$ positive and invertible (for any x),
- ▶ u, p^μ, q have homogeneous gauge and diff. transformations (sections of $\text{End}(V)$),
- ▶ $\Gamma^\nu := g^{\mu\rho} \Gamma_{\mu\rho}^\nu$ where the $\Gamma_{\mu\rho}^\nu =$ Christoffel symbols of g .
- For any $a \in \Gamma(\text{End}(V))$, consider the asymptotics of the heat-trace

$$\text{Tr}[ae^{-tP}] \underset{t \downarrow 0^+}{\sim} \sum_{r=0}^{\infty} a_r(a, P) t^{(r-d)/2}$$

- **Compute \mathcal{R}_r defined by**

$$a_r(a, P) = \int_M \text{tr}[a(x)\mathcal{R}_r(x)] \text{dvol}_g(x).$$

The key steps for the computation

Iochum, B. and Masson, T. (2017). Heat trace for Laplace type operators with non-scalar symbols. *Journal of Geometry and Physics* 116, pp. 90–118

Iochum, B. and Masson, T. (2018). Heat asymptotics for nonminimal Laplace type operators and application to noncommutative tori. *Journal of Geometry and Physics* 129, pp. 1–24

- $\text{Tr}[ae^{-tP}] = \int dx \text{tr}[a(x)K(t, x, x)]$ with $K(t, x, x)$ diagonal of the kernel of e^{-tP} :

$$\int_M dy K(t, x, y)s(y) = (e^{-tP}s)(x), \text{ for any } s \in \Gamma(V)$$

- It is well known that

$$K(t, x, y) = (2\pi)^{-d} \int_{\mathbb{R}^d} d\xi e^{-iy \cdot \xi} (e^{-tP} e^{ix \cdot \xi}) : \mathbb{C}^N \rightarrow \mathbb{C}^N.$$

- One has (first trick):

$$-Pe^{ix\xi}s = -e^{ix\xi}[H + K + P]s, \text{ with } \begin{cases} K := -i\xi_\mu(p^\mu + g^{\mu\nu}(\nabla_\nu u) - \Gamma^\mu u + 2g^{\mu\nu}u\nabla_\nu) \\ H := g^{\mu\nu}u\xi_\mu\xi_\nu \end{cases}$$

then, at fixed x , $K(t, x, x)$ is the linear operator (on \mathbb{C}^N):

$$\mathbb{C}^N \ni v \mapsto K(t, x, x)v = t^{-d/2}(2\pi)^{-d} \int_{\mathbb{R}^d} d\xi e^{-H - \sqrt{t}K - tP} v.$$

The key steps (cont'd)

- One uses (second trick) the Volterra series

$$e^{A+B} = e^A + \sum_{k=1}^{\infty} \int_{\Delta_k} ds e^{(1-s_1)A} B e^{(s_1-s_2)A} \dots e^{(s_{k-1}-s_k)A} B e^{s_k A},$$

where $\Delta_k := \{s = (s_1, \dots, s_k) \in \mathbb{R}_+^k \mid 0 \leq s_k \leq s_{k-1} \leq \dots \leq s_2 \leq s_1 \leq 1\}$ ($\Delta_0 := \emptyset$)

- Then

$$e^{-H-\sqrt{t}K-tP} v = e^{-H} v + \sum_{k=1}^{\infty} (-1)^k f_k(\xi) [(\sqrt{t}K + tP) \otimes \dots \otimes (\sqrt{t}K + tP)] v$$

where, for any $k \in \mathbb{N}$, the map $f_k(\xi) : M_N[\xi, \nabla]^{\otimes k} \rightarrow M_N[\xi]$ is defined by

$$f_k(\xi)[B_1 \otimes \dots \otimes B_k] := \int_{\Delta_k} ds e^{(s_1-1)H} B_1 e^{(s_2-s_1)H} B_2 e^{(s_3-s_2)H} \dots B_k e^{-s_k H},$$

$$f_0(\xi)[z] := z e^{-H}, \quad \text{for } z \in \mathbb{C} =: M_N^{\otimes 0}.$$

- The B_i are matrix-valued diff. op. in ∇_μ depending on x and (polynomials in) ξ .

The key steps (cont'd)

- Then one has to compute:

$$a_0(a, P)(x) = \operatorname{tr} \frac{|g|^{-1/2}}{(2\pi)^d} a(x) \int_{\mathbb{R}^d} d\xi f_0[1],$$

$$a_2(a, P)(x) = \operatorname{tr} \frac{|g|^{-1/2}}{(2\pi)^d} a(x) \int_{\mathbb{R}^d} d\xi (f_2[K \otimes K] - f_1[P]),$$

$$a_4(a, P)(x) = \operatorname{tr} \frac{|g|^{-1/2}}{(2\pi)^d} a(x) \int_{\mathbb{R}^d} d\xi (f_2[P \otimes P] - f_3[K \otimes K \otimes P] - f_3[K \otimes P \otimes K] - f_3[P \otimes K \otimes K] + f_4[K \otimes K \otimes K \otimes K]),$$

etc.

- Strategy for the computation:

- 1 Compute and characterize the operators $f_k : M_N[\xi, \nabla]^{\otimes k} \rightarrow M_N[\xi]$;
- 2 Take care of the operators ∇_μ in the arguments...

Let us use the notations:

$$\begin{aligned} H^{\mu\nu} &:= g^{\mu\nu} u, & L^\mu &:= p^\mu + g^{\mu\nu}(\nabla_\nu u) - \Gamma^\mu u, & K^\mu &:= -i(L^\mu + 2H^{\mu\nu}\nabla_\nu), \\ H &= H^{\mu\nu}\xi_\mu\xi_\nu, & K &= K^\mu\xi_\mu, & P &= -H^{\mu\nu}\nabla_\mu\nabla_\nu - L^\mu\nabla_\mu - q. \end{aligned}$$

The universal operators

- Define the functions of $r_i \in \text{spectrum of } u(x)$ (x is fixed):

$$I_{\alpha,k}(r_0, r_1, \dots, r_k) := \int_{\Delta_k} ds [(1-s_1)r_0 + (s_1-s_2)r_1 + \dots + s_k r_k]^{-\alpha}$$

- Define the associated (x -dependent) operators $X_{\alpha,k} \in \mathcal{B}(M_N^{\otimes k}, M_N)$ by

$$X_{\alpha,k}[B_1 \otimes \dots \otimes B_k] := I_{\alpha,k}(r_0, \dots, r_k) E_0 B_1 E_1 \dots B_k E_k$$

with summation over $k+1$ -uplets (r_0, \dots, r_k) of spectral values of $u(x)$, E_i the spectral projection of $u(x)$ associated to r_i .

- One has

$$X_{d/2+p,k,\mu_1 \dots \mu_{2p}} := G_{\mu_1 \dots \mu_{2p}} X_{d/2+p,k} = \frac{1}{g_d} \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} d\xi \xi_{\mu_1} \dots \xi_{\mu_{2p}} f_k(\xi)$$

with

$$G_{\mu_1 \dots \mu_{2p}} := \frac{1}{2^{2p} p!} \sum_{\rho \in S_{2p}} g_{\mu_{\rho(1)} \mu_{\rho(2)}} \dots g_{\mu_{\rho(2p-1)} \mu_{\rho(2p)}}$$

(S_{2p} = symm. group of perm. on $2p$ elements) and $g_d = \frac{|g|^{1/2}}{2^d \pi^{d/2}}$.

- Everything here depend only on u and $g^{\mu\nu}$, not on P .**

Propagation of the ∇_μ

- To compute \mathcal{R}_r , we start with terms of the form

$$\begin{aligned} \frac{1}{(2\pi)^d} \int d\xi \xi^{\mu_1} \cdots \xi^{\mu_{2p}} f_k(\xi) [(B_1 \otimes \cdots \otimes B_k)^{\mu_1 \cdots \mu_{2p}}] \\ = g_d X_{d/2+p, k, \mu_1 \cdots \mu_{2p}} [(B_1 \otimes \cdots \otimes B_k)^{\mu_1 \cdots \mu_{2p}}] \end{aligned}$$

- Let $Q[A]$ be a matrix-valued function written as a polynomial in the A_μ and their derivatives (∇_μ repeatedly acting on $v \in \mathbb{C}^N$). One main result is

$$\begin{aligned} X_{d/2+p, k, \mu_1 \cdots \mu_{2p}} [(B_1 \otimes \cdots \otimes B_i \nabla_v \otimes \cdots \otimes B_k)^{\mu_1 \cdots \mu_{2p}}] Q[A] \\ = \sum_{j=i+1}^k X_{d/2+p, k, \mu_1 \cdots \mu_{2p}} [(B_1 \otimes \cdots \otimes (\nabla_v B_j) \otimes \cdots \otimes B_k)^{\mu_1 \cdots \mu_{2p}}] Q[A] \\ - \sum_{j=i}^k X_{d/2+p+1, k+1, \mu_1 \cdots \mu_{2(p+1)}} [(B_1 \otimes \cdots \otimes B_j \otimes (\nabla_v H^{\mu_{2p+1} \mu_{2p+2}}) \otimes \cdots \otimes B_k)^{\mu_1 \cdots \mu_{2p}}] Q[A] \\ + X_{d/2+p, k, \mu_1 \cdots \mu_{2p}} [(B_1 \otimes \cdots \otimes B_i \otimes \cdots \otimes B_k)^{\mu_1 \cdots \mu_{2p}}] (\nabla_v Q[A]). \end{aligned}$$

Use repeatedly \rightarrow all the ∇_v are applied...

Computation of \mathcal{R}_2

$$H^{\mu\nu} := g^{\mu\nu} u,$$

$$H = H^{\mu\nu} \xi_\mu \xi_\nu,$$

$$L^\mu := p^\mu + g^{\mu\nu} (\nabla_\nu u) - \Gamma^\mu u,$$

$$K = K^\mu \xi_\mu,$$

$$K^\mu := -i(L^\mu + 2H^{\mu\nu} \nabla_\nu),$$

$$P = -H^{\mu\nu} \nabla_\mu \nabla_\nu - L^\mu \nabla_\mu - q.$$

- Start with

$$\mathcal{R}_2 = \frac{1}{2^d \pi^{d/2}} \left(X_{d/2+1,2,\mu_1\mu_2} [K^{\mu_1} \otimes K^{\mu_2}] - X_{d/2,1} [P] \right).$$

- Propagate the $\nabla_\mu \rightarrow$ **24 terms** (after simplifications).
- Replace $H^{\mu\nu}$, L^μ and their derivatives (use normal coordinates)...

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- Propagate the $\nabla_\mu \rightarrow$ **24 terms** (after simplifications).
- Replace $H^{\mu\nu}$, L^μ and their derivatives (use normal coordinates)...
- Final result, with $\widehat{\nabla}_\mu$ the “full covariant derivative” (gauge + Riemannian):

$$\begin{aligned} \mathcal{R}_2 = & \frac{1}{2^d \pi^{d/2}} \left(\frac{1}{6} \mathfrak{R} X_{d/2,1}[u] + X_{d/2,1}[q] - X_{d/2+1,2}[u \otimes \widehat{\nabla}_\mu p^\mu] \right. \\ & + \frac{1}{2} g^{\mu\nu} X_{d/2+1,2}[(\widehat{\nabla}_\mu u + p_\mu) \otimes (\widehat{\nabla}_\nu u - p_\nu)] \\ & - \frac{d+2}{2} g^{\mu\nu} X_{d/2+2,3}[u \otimes \widehat{\nabla}_{\mu\nu}^2 u \otimes u] \\ & - \frac{d+2}{2} g^{\mu\nu} X_{d/2+2,3}[u \otimes \widehat{\nabla}_\mu u \otimes (\widehat{\nabla}_\nu u - p_\nu)] \\ & - \frac{d+2}{2} g^{\mu\nu} X_{d/2+2,3}[(\widehat{\nabla}_\mu u + p_\mu) \otimes \widehat{\nabla}_\nu u \otimes u] \\ & \left. + \frac{(d+2)(d+4)}{2} g^{\mu\nu} X_{d/2+3,4}[u \otimes \widehat{\nabla}_\mu u \otimes \widehat{\nabla}_\nu u \otimes u] \right). \end{aligned}$$



Computation of \mathcal{R}_4

- Start with

$$\begin{aligned} \mathcal{R}_4 = \frac{1}{2^d \pi^{d/2}} & \left(X_{d/2,2} [P \otimes P] - X_{d/2+1,3,\mu_1\mu_2} [K^{\mu_1} \otimes K^{\mu_2} \otimes P] \right. \\ & - X_{d/2+1,3,\mu_1\mu_2} [K^{\mu_1} \otimes P \otimes K^{\mu_2}] - X_{d/2+1,3,\mu_1\mu_2} [P \otimes K^{\mu_1} \otimes K^{\mu_2}] \\ & \left. + X_{d/2+2,4,\mu_1\mu_2\mu_3\mu_4} [K^{\mu_1} \otimes K^{\mu_2} \otimes K^{\mu_3} \otimes K^{\mu_4}] \right). \end{aligned}$$

- Propagate the $\nabla_\mu \rightarrow$ **3771 terms** (after simplifications!).

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\mathcal{R}_4 : computer enters the game...

- All the manipulations in the method are “symbolic” and “combinatorial”.
- A **computer algebra system** should do the work for us!
- Requirements:
 - ▶ Derivations (∇_μ);
 - ▶ Noncommutative products (matrix algebra)...
 - ▶ mixed with commutative products (metric...);
 - ▶ Tensor products;
 - ▶ Riemannian tensor manipulations (in normal coordinates);
 - ▶ Gauge field strength and its derivatives;
 - ▶ Contractions of indices;
 - ▶ Simplification of expressions combining all these structures...

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 - ▶ Simplification of expressions combining all these structures...
- First tentative: Mathematica...
 - ➔ Not easy to make it understand all these requirements!
- Present tentative: **home made object oriented software**.
 - ▶ We manipulate (programmed) “objects” that reproduce our (math.) “objects”...
 - ▶ Can be written in lot of languages: we (I) chose Javascript with Node.
(I am used to Javascript, favorable benchmarks compared to Python...)



Problems and benefits...

- 👎 We (Bruno and me) are not specialist in “formal mathematical manipulations” on computers...
 - ▶ How to simplify?
 - ▶ How to tell the software that we want a expression instead of another?
- 👎 How to be sure that the software computes what we want?
- 👍 The software is designed for the method and we can add filters, post operations (simplifications), etc...
- 👍 \LaTeX exportation...

Some preliminary results...

- The 3771 terms can be sorted according to the only 5 values of $Q[A]$: $\mathbb{1}$ (2026 terms), $\nabla_{v_1} \mathbb{1}$ (1296 terms), $\nabla_{v_1 v_2}^2 \mathbb{1}$ (382 terms), $\nabla_{v_1 v_2 v_3}^3 \mathbb{1}$ (62 terms), and $\nabla_{v_1 v_2 v_3 v_4}^4 \mathbb{1}$ (5 terms).
- $Q[A] = \nabla_{v_1} \mathbb{1}$: no gauge homogeneous expression \rightarrow should be 0.
This is what the computer returns!
- $Q[A] = \nabla_{v_1 v_2 v_3 v_4}^4 \mathbb{1}$: the only gauge homogeneous expression is $F_{\mu\nu} F^{\mu\nu}$.
The computer produces (directly) only one term:

$$\frac{1}{12} X_1[u] F^{v_1 v_2} F_{v_1 v_2}.$$

Special case $u = \mathbb{1}$: this reduces to $\frac{1}{12} F^{v_1 v_2} F_{v_1 v_2}$ and agrees with results in [Gilkey, P. B. \(2003\). *Asymptotic formulae in spectral geometry*. CRC press.](#)

Some preliminary results (cont'd)...

- $Q[A] = \nabla_{v_1 v_2 v_3}^3 \mathbf{1}$: let $\mathcal{X}_{(3)}[a] := \frac{1}{6} X_1[a] - X_3[u \otimes a \otimes u]$, then

$$\begin{aligned} & - 4 g^{v_1 v_2} g^{v_3 v_4} X_5[(\widehat{\nabla}_{v_4} u) \otimes u \otimes u \otimes u \otimes u] (\widehat{\nabla}_{v_1} F_{v_2 v_3}) \\ & + d g^{v_1 v_2} g^{v_3 v_4} X_5[u \otimes (\widehat{\nabla}_{v_4} u) \otimes u \otimes u \otimes u] (\widehat{\nabla}_{v_1} F_{v_2 v_3}) \\ & - d g^{v_1 v_2} g^{v_3 v_4} X_5[u \otimes u \otimes u \otimes (\widehat{\nabla}_{v_4} u) \otimes u] (\widehat{\nabla}_{v_1} F_{v_2 v_3}) \\ & + 4 g^{v_1 v_2} g^{v_3 v_4} X_5[u \otimes u \otimes u \otimes u \otimes (\widehat{\nabla}_{v_4} u)] (\widehat{\nabla}_{v_1} F_{v_2 v_3}) \\ & - g^{v_1 v_2} \mathcal{X}_{(3)}[p^{v_3}] (\widehat{\nabla}_{v_1} F_{v_2 v_3}). \end{aligned}$$

➔ more simplifications to expect?

- $Q[A] = \nabla_{v_1 v_2}^2 \mathbf{1}$: work in progress, we analyze what the software has returned. “Post-production” on 31 terms...
- $Q[A] = \mathbf{1}$: that's the big piece of the cake!
 - ▶ Use normal coordinates up to 4-th derivatives of $H^{\mu\nu}$...
 - ▶ Replace gauge covariant derivatives with full covariant derivatives...
 - ▶ Introduce derivations of Riemann and Ricci tensors...
 - ▶ Self (full) contractions of Riemann tensor and Ricci tensor...
 - ▶ Lot of work to be done “in post production” to get (human) readable expression...

Conclusions

- The method gives “universal” formulas in terms of universal operators.
- These formulas can be instantiated to specific situations:
 - ▶ “diagonal” case $u = \mathbb{1} \Rightarrow$ know results...
 - ▶ Noncommutative Torus \Rightarrow we confirm results by Connes, Tretkoff, Moscovici, Fathizadeh, Khalkhali... (avoid NC pseudodifferential calculus).
- Search for “hidden” structures (patterns) in the result.
Wait and see what \mathcal{R}_4 will tell us...
- The software could help compute \mathcal{R}_r for large r ...
with the help of a powerful computer?
- The method could (should!) work in NCG:
spectral decompositions, algebraic constructions, traces...
- Only work with leading term $-g^{\mu\nu} u \nabla_\mu \nabla_\nu$ in P ...